

Scattering theory and operator algebras: a fruitful exchange

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Scattering theory: a comparison theory

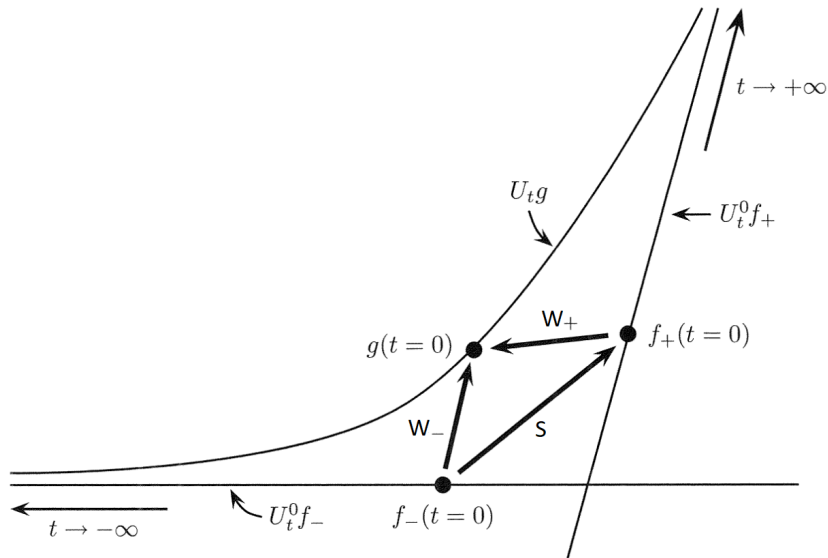


Figure: Representation of scattering theory

Scattering theory (self-adjoint theory)

$\mathcal{H}$ a Hilbert space, H_0, H self-adjoint operators in $\mathcal{H}$ with

- H_0 purely absolutely continuous,
- $W_{\pm} := s - \lim_{t \rightarrow \pm\infty} e^{itH} e^{-itH_0} E_{ac}(H_0)$ exist,
- $\text{Ran}(W_-) = \text{Ran}(W_+) = (1 - E_p(H))\mathcal{H}$.

Then,

- The **wave operators** $W_{\pm}$ are isometries with

$$W_{\pm}^* W_{\pm} = 1 \quad \text{and} \quad W_{\pm} W_{\pm}^* = 1 - E_p(H),$$

- The **scattering operator** $S := W_+^* W_-$ is unitary.

Index map in K -theory [RLL]

Let

$$0 \longrightarrow \mathcal{J} \hookrightarrow \mathcal{E} \xrightarrow{q} \mathcal{Q} \longrightarrow 0$$

be a short sequence of C^* -algebras, with $\mathcal{E}$ unital. Let U be a unitary element in $\mathcal{Q}$ which has a lift to a partial isometry W in $\mathcal{E}$. Then $1 - W^*W$ and $1 - WW^*$ are projections in $\mathcal{J}$, and

$$\operatorname{ind}([U]_1) = [1 - W^*W]_0 - [1 - WW^*]_0.$$

If $W_- \in \mathcal{E}$ with $q(W_-)$ unitary, then

$$\operatorname{ind}([q(W_-)]_1) = -[E_p(H)]_0.$$

But how to choose $\mathcal{E}$? Who is $q(W_-)$? Can we get an equality between numbers ?

Filling the framework: Topological Levinson's Thm

The program:

- Fix a pair (H, H_0) ,
- Get a rather explicit formula for W_- ,
- Guess a suitable algebra $\mathcal{E}$ with $W_- \in \mathcal{E}$,
- Find $\mathcal{J}$ such that $q(W_-)$ is unitary,

We then get

$$\text{ind} \left([q(W_-)]_1 \right) = -[E_p(H)]_0.$$

- Deduce a numerical equality from this relation.

This program has been successful on several examples.

¿ What can scattering theory bring back to operator algebras ?

Example I: The operators $H_{m,\kappa}$

For $|\operatorname{Re}(m)| < 1$, $m \neq 0$, and $\kappa \in \mathbb{C} \cup \{\infty\}$ we set in $L^2(\mathbb{R}_+)$

$$H_{m,\kappa} := -\frac{d^2}{dr^2} + \left(m^2 - \frac{1}{4}\right) \frac{1}{r^2}$$

with

$$\mathcal{D}(H_{m,\kappa}) := \left\{ f \in L^2(\mathbb{R}_+) \mid H_{m,\kappa} f \in L^2(\mathbb{R}_+) \text{ and } \exists c \in \mathbb{C} \text{ s.t.} \right. \\ \left. f(r) \sim c(\kappa r^{\frac{1}{2}-m} + r^{\frac{1}{2}+m}) \text{ near } 0 \right\}.$$

$H_{m,\kappa}$ is self-adjoint if (m, κ) belongs to

$$\Omega_{\text{sa}} := \{(m, \kappa) \mid m \in (-1, 1) \setminus \{0\} \text{ and } \kappa \in \mathbb{R}, \\ \text{or } m \in i\mathbb{R} \setminus \{0\} \text{ and } |\kappa| = 1\},$$

and $\#\sigma_p(H_{m,\kappa}) = \dim(E_p(H_{m,\kappa})) \in \mathbb{N} \cup \{\infty\}$.

Example I: The wave operators

$W_-(H_{m,\kappa}, H_{m',\kappa'})$ are partial isometries and unitarily equivalent in $L^2(\mathbb{R})$ to

$$\begin{aligned} & \mathcal{W}_-(H_{m,\kappa}, H_{m',\kappa'}) \\ &= \Xi_{\frac{1}{2}}(-D) \left(\Xi_m(D) - \varsigma \Xi_{-m}(D) e^{2mX} \right) \frac{e^{-i\frac{\pi}{2}m}}{1 - \varsigma e^{-i\pi m} e^{2mX}} \\ & \quad \times \frac{e^{+i\frac{\pi}{2}m'}}{1 - \varsigma' e^{+i\pi m'} e^{2m'X}} \left(\Xi_{m'}(-D) - \varsigma' e^{2m'X} \Xi_{-m'}(-D) \right) \Xi_{\frac{1}{2}}(D). \end{aligned}$$

with $\varsigma := \kappa \frac{\Gamma(-m)}{\Gamma(m)}$, $\varsigma' := \kappa' \frac{\Gamma(-m')}{\Gamma(m')}$ and

$$\mathbb{R} \ni \xi \mapsto \Xi_m(\xi) := e^{i \ln(2)\xi} \frac{\Gamma\left(\frac{m+1+i\xi}{2}\right)}{\Gamma\left(\frac{m+1-i\xi}{2}\right)} \in \mathbb{C}.$$

Example I: Classification

$\mathcal{W}_-(H_{m,\kappa}, H_{m',\kappa'})(x, \xi)$ is a Ψ DO of order 0 in ξ , and...

The Fredholm case

- The coef. have limits at $x = \pm\infty$,
- Already existing framework, classical situation, studied in [NPR].

The semi-Fredholm case (I)

- The coef. are periodic,
- Atiyah's L^2 -index theorem [Ati].

The semi-Fredholm case (II)

- The coef. are asymptotically periodic,
- C^* -algebras studied in [CM],
- 2-link ideal chain and two different index theorems.

The almost periodic case

- The coef. are almost periodic,
- C^* -algebras and index theorems in [CDSS, CMS, Shu].

Example I: The Fredholm case

Let $m \in (-1, 1) \setminus \{0\}$, $\kappa \in \mathbb{R}$, and $(m', \kappa') = (\frac{1}{2}, 0)$. Then,

$$\mathcal{W}_-(H_{m,\kappa}, H_{\frac{1}{2},0}) \in \mathcal{E}_\square := C^*\left(a(X)b(D) \mid a, b \in C([-\infty, \infty])\right).$$

- $\mathcal{E}_\square / \mathcal{K}(L^2(\mathbb{R})) \cong C(\square) \cong C(\mathbb{S}^1)$
- The quotient map q is given by

$$a(X)b(D) \mapsto \left(b(-\infty)a, a(+\infty)b, b(+\infty)a, a(-\infty)b\right)$$

- $\Gamma_{m,\kappa;\frac{1}{2},0} := q(\mathcal{W}_-(H_{m,\kappa}, H_{\frac{1}{2},0}))$ has 4 components, one of them is unitarily equivalent to the scattering operator.

Theorem (Fredholm case)

$$\text{Wind}_\square(\Gamma_{m,\kappa;\frac{1}{2},0}) = -\text{Index}(\mathcal{W}_-(H_{m,\kappa}, H_{\frac{1}{2},0})) = \#\sigma_p(H_{m,\kappa}).$$

Example II: Unitary theory (discrete time)

In $\mathcal{H} := \ell^2(\mathbb{Z} \times \mathbb{N})$, let U_0, U_1 be unitary operators,

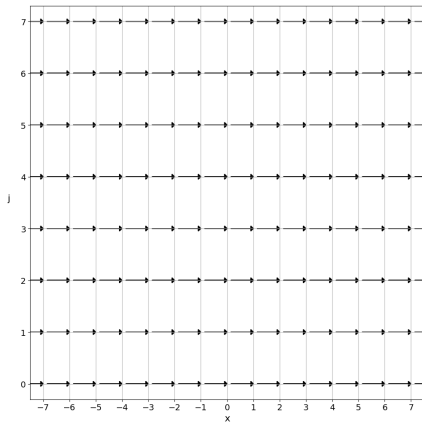


Figure: Action of U_0

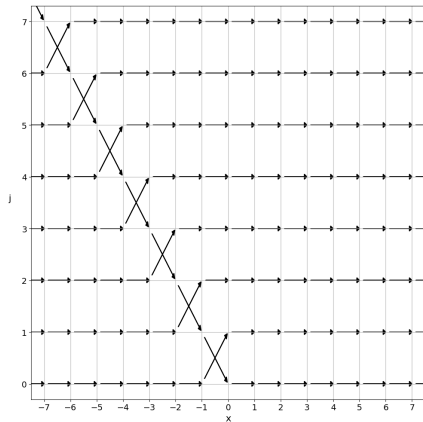


Figure: Action of U_1

and the wave operators $W_{\pm}(U_1, U_0) := s - \lim_{n \rightarrow \pm\infty} U_1^{-n} U_0^n$.

Example II: Wave operators $W_{\pm}(U_1, U_0)$

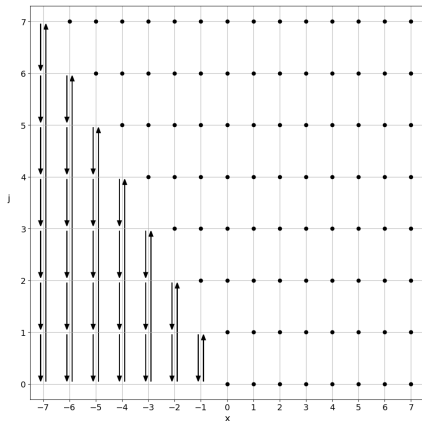


Figure: Action of $W_+(U_1, U_0)$

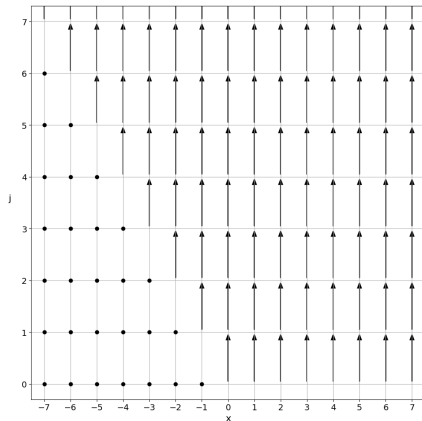


Figure: Action of $W_-(U_1, U_0)$

Observe that $\text{Ran}(W_+(U_1, U_0)) = \mathcal{H}$ but $\text{Ran}(W_-(U_1, U_0)) \neq \mathcal{H}$.

Example II: U_2 with one invariant point

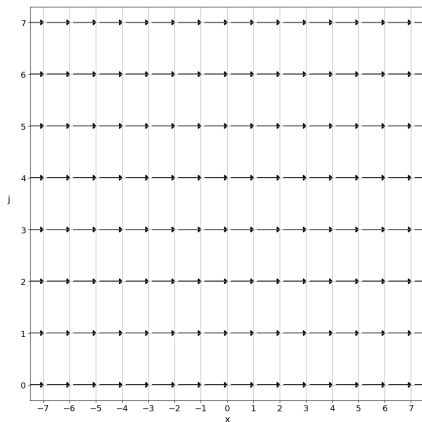


Figure: Action of U_0

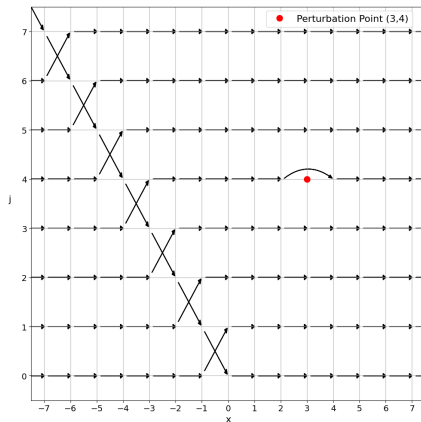


Figure: Action of U_2

and the wave operators $W_{\pm}(U_2, U_0) := s - \lim_{n \rightarrow \pm\infty} U_2^{-n} U_0^n$.

Example II: Wave operators $W_{\pm}(U_2, U_0)$

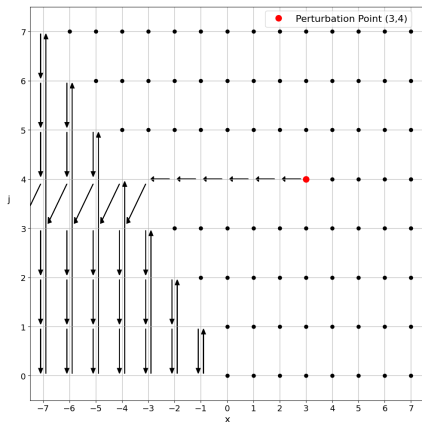


Figure: Action of $W_+(U_2, U_0)$

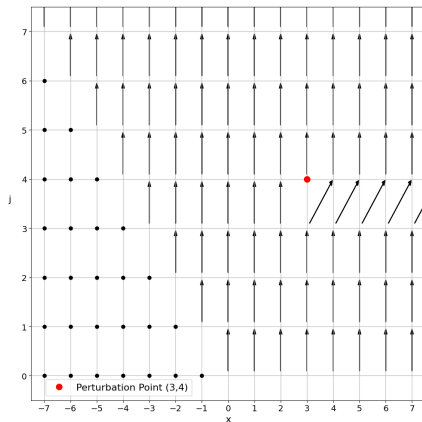


Figure: Action of $W_-(U_2, U_0)$

¿ All these operators are isometries, what do they illustrate ?

Example II: Wold's decomposition

On $\ell^2(\mathbb{N})$ with the standard basis $\{\delta_n\}_{n \in \mathbb{N}}$, set $S\delta_n = \delta_{n+1}$.

Theorem (Wold's decomposition)

If W is an isometry on a Hilbert space $\mathcal{H}$, then there is a cardinal number α and a unitary operator V (possibly vacuous) such that W is unitarily equivalent to $S^{(\alpha)} \oplus V$.

These examples provide a perfect visualization of this theorem. Other general properties of scattering theory can be illustrated with these operators.

Conclusion

- Operator algebras provides a natural framework for new investigations on scattering theory
- In return scattering theory provides non-trivial examples
- What is the next step ?

THANK YOU

References

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